

Lecture #4 of 12

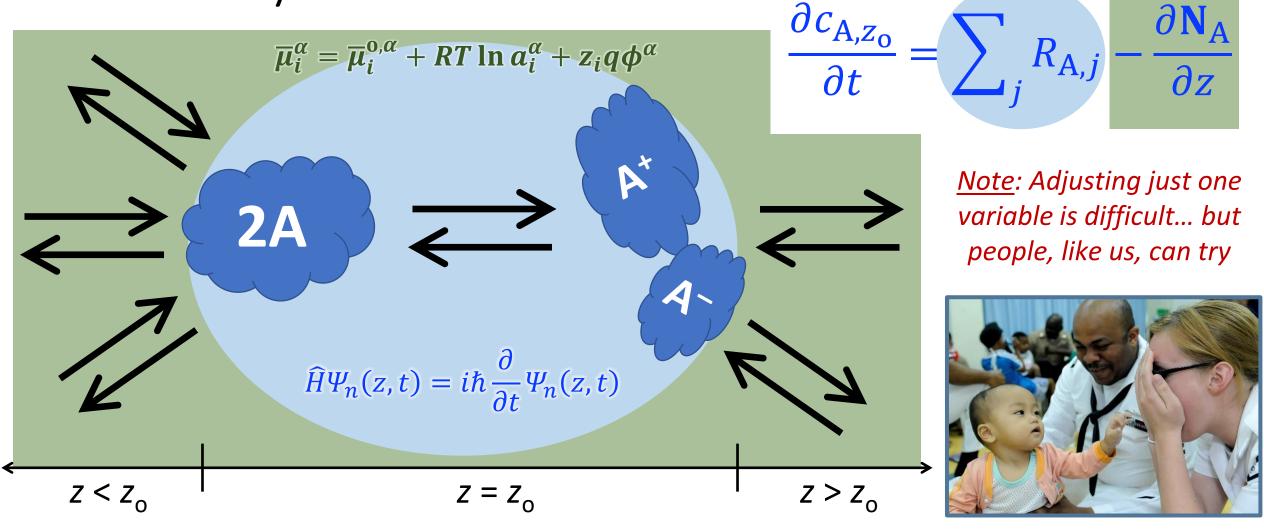
Prof. Shane Ardo
Department of Chemistry
University of California Irvine



Thermal (Dark) Reactions

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Continuity of Mass



"Jean Piaget, the Swiss psychologist who first studied <u>object permanence</u> in infants, argued that it is one of an infant's most important accomplishments, as, without this concept, objects would have no separate, permanent existence."

— Wiki, Object permanence

Transport for Photochemistry

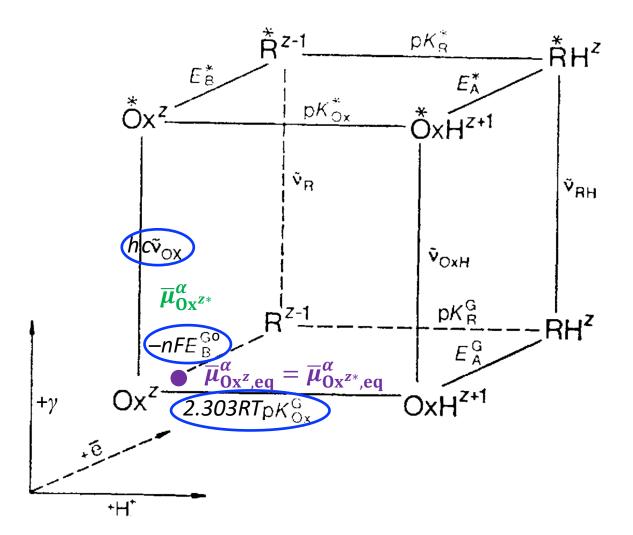
So, what are ideal mass-transfer-related properties in photochemical systems?

$$\operatorname{Flux}_{z,e} = \mathbf{N}_e = -\frac{\sigma_e}{F^2} \frac{d\boldsymbol{\mu}_e}{dz}$$

- (1) Transport is fast = flux is large...
 - ... so that excited states can always interact with reactants to perform chemical reactions ... so that reactants are not depleted by chemical reactions
- (2) At least 2 non-equilibrium species transport to different locations = "selective contacts"...
 - ... which results in different observed rates for various chemical reactions at each location
 - ... that are ideally oppositely charged, so that transport of only a few results in large Φ , $\Delta \overline{\mu}_i$
- (2') Reactivity by non-equilibrium species is selective = desired reactions are fast...
 - ... so that photon energy can be stored in chemical products, when $\Delta G>0$ for the reaction
 - ... so that photon energy can be used to speed up reactions with $\Delta G < 0$

$$\operatorname{Flux}_{z,e} = -D_e n_e \left(\frac{d\mu_e^0}{kT_e dz} + \frac{d(\ln \gamma_e)}{dz} + \frac{d(\ln n_e)}{dz} - \frac{d(\ln n_e^0)}{dz} - \frac{d(\ln n_e^0)}{dz} - \frac{d(q\varphi)}{kT_e dz} \right) - D_{n_e,\gamma_e,n_e^0,T_e} \frac{dT_e}{dz}$$

Förster Cube and Square Schemes



... all of these free energy <u>terms</u> are **standardstate** free energies (ΔG°)... but what is the actual free energy of the system (ΔG)?

... let's assume that $\Delta G = 0$ (equilibrium)... how could I indicate that on this slide, as a point(s), to depict the majority species present?

... now, how can one **push/pull** this system out of equilibrium?

... recall Le Châtelier's principle... and thus by addition of reactants or removal of products... such as mass <u>or light!</u>

... hopefully this made a little more sense this time around... and if not, let's keep on trying!

Summary of Key Equations and Equilibrium (UPDATED) 79

 $\frac{\partial c_{A,Z_0}}{\partial t} = \sum_{j} R_{A,j}$

(Chemists should be yawning)

$$\overline{\mu}_{i}^{\alpha} = \left(\frac{\partial G^{\alpha}}{\partial n_{i}^{\alpha}}\right)_{T,n,n_{i+1}^{\alpha}}$$

 $A \rightleftharpoons B$

$$\overline{\mu}_{B}^{\alpha} - \overline{\mu}_{A}^{\alpha} = \Delta G^{\alpha} = \Delta G^{0,\alpha} + RT \ln Q$$
... where $Q = \frac{a_{B}^{\alpha}}{a_{A}^{\alpha}} = \frac{\gamma_{B}^{\alpha}(c_{B}^{\alpha}/c_{B}^{0,\alpha})}{\gamma_{A}^{\alpha}(c_{A}^{\alpha}/c_{A}^{0,\alpha})}$

$$R_{\mathrm{A},f} = -\mathbf{k_f} a_{\mathrm{A}}$$
 $R_{\mathrm{B},f} = +\mathbf{k_f} a_{\mathrm{A}}$ $\mathbf{k_j}$ (M/s)... a rate!

$$R_{\mathrm{A},b} = + k_b a_{\mathrm{B}}$$
 $R_{\mathrm{B},b} = -k_b a_{\mathrm{B}}$ k_j' (s⁻¹)... an inverse time constant!

$$\overline{\mu}_{i}^{\alpha} = \left(\frac{\partial G^{\alpha}}{\partial n_{i}^{\alpha}}\right)_{T n n^{\alpha}} \qquad \frac{\partial c_{A,Z_{0}}^{\alpha}}{\partial t} = -k_{f} a_{A}^{\alpha} + k_{b} a_{B}^{\alpha} = -k_{f}^{\prime} c_{A}^{\alpha} + k_{b}^{\prime} c_{B}^{\alpha} = -\frac{\partial c_{B,Z_{0}}^{\alpha}}{\partial t}$$

Now, consider the reaction to be at equilibrium (3 ways)... $\Delta G^{\alpha} = 0 = \Delta G^{0,\alpha} + RT \ln Q$

... and so,
$$\Delta G^{\text{o},\alpha} = -RT \ln K$$
, with $K = \frac{a_{\text{B,eq}}^{\alpha}}{a_{\text{A,eq}}^{\alpha}} = \frac{\gamma_{\text{B}}^{\alpha} \left(c_{\text{B,eq}}^{\alpha} / c_{\text{B}}^{\text{o},\alpha} \right)}{\gamma_{\text{A}}^{\alpha} \left(c_{\text{A,eq}}^{\alpha} / c_{\text{A}}^{\text{o},\alpha} \right)}$

And also, $\bar{\mu}_{A}^{\alpha} - \bar{\mu}_{A}^{\alpha} = 0$... and so, $\bar{\mu}_{A}^{\alpha} = \bar{\mu}_{B}^{\alpha}$... if that is even helpful

And also,
$$\frac{\partial c_{A,z_0}^{\alpha}}{\partial t} = \mathbf{0} = -\frac{\partial c_{B,z_0}^{\alpha}}{\partial t} = -\mathbf{k_f} a_{A,eq}^{\alpha} + \mathbf{k_b} a_{B,eq}^{\alpha}$$
 and so $\frac{\mathbf{k_f}}{\mathbf{k_b}} = \frac{a_{B,eq}^{\alpha}}{a_{A,eq}^{\alpha}} = \frac{\gamma_{B}^{\alpha} \left(c_{B,eq}^{\alpha} / c_{B}^{o,\alpha} \right)}{\gamma_{A}^{\alpha} \left(c_{A,eq}^{\alpha} / c_{A}^{o,\alpha} \right)} = K$

Today's Critical Guiding Question

What thermodynamic term indicates whether a processes will net occur, or not, for a chemical reaction, and a single species... and uniquely for today, how does it allow us to quantitatively predict kinetic information... for real this time?

Thermal (Dark) Reactions

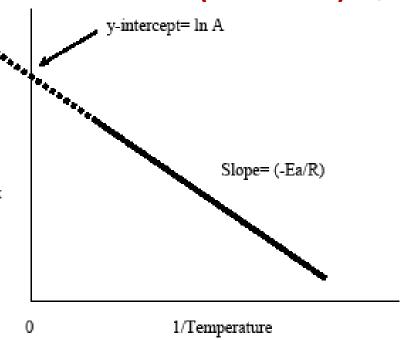
- Activation energy, Eyring—Polanyi—Evans equation
- Marcus-Hush (electron-transfer) theory
- Transition-state character, Reorganization energies (outer and inner), Linear free energy relationships
- Molecular orbital theory, Huang–Rhys factor
- Quantum mechanical tunneling, Superexchange
- Inner versus Outer sphere reactions, Robin—Day classification
- Self-exchange reactions, Marcus cross relations
- Charge transfer across electrified interfaces, Butler–Volmer equation,
 Rate-determining step, Fermi's golden rule, Marcus–Gerischer theory

Activation Energy

Arrhenius (1889)

empirical rate constant equation

$$k_f = Ae^{-\frac{E_a}{RT}}$$
 ... pre-exponential factor, A , has units of s⁻¹ ln ln



$A \rightleftharpoons X^{\neq} \rightleftharpoons B$

Eyring-Polanyi-Evans (1930s)

theoretical rate constant equation (from transition-state theory / activated complex theory)

$$k_f = \kappa \nu K^{\neq} = \frac{\kappa k_{\rm B} T}{h} \frac{K^{\neq}}{\text{this } N_{\rm A}} \text{ is the Avogadro constant ... and } R \stackrel{\neq}{=} N_A k_{\rm B} ... \text{ and } \frac{\kappa k_{\rm B} T}{h} \text{ has units of s}^{-1}$$

 $A \rightleftharpoons B$

$$k_f = \frac{\kappa k_{\rm B} T}{h} \exp\left(-\frac{\Delta G^{\neq}}{RT}\right) = \frac{\kappa k_{\rm B} T}{h} \exp\left(\frac{\Delta S^{\neq}}{R}\right) \exp\left(-\frac{\Delta H^{\neq}}{RT}\right) \text{ ... and so } A \text{ contains } \Delta S^{\neq}$$

... what is the largest predicted pre-exponential factor at 25 °C? (161 fs)⁻¹ = (1.61 x 10⁻¹³ s)⁻¹

Marcus-Hush (1950s-1960s)

theoretical (semiclassical) rate constant equation

 $\frac{D + A}{A} \rightleftharpoons \frac{D^{+} + A^{-}}{B}$... but how is this a first-order reaction (s⁻¹)?... more on this later...

$$k_{\mathrm{ET}} = rac{2\pi}{\hbar} |H_{\mathrm{DA}}|^2 rac{1}{\sqrt{4\pi\lambda_{\mathrm{AB}}k_{\mathrm{B}}T}} \exp\left(-rac{\Delta G_{\mathrm{AB}}^{\ \ \ \ }}{k_{\mathrm{B}}T}
ight) \ ... rac{2\pi |H_{\mathrm{DA}}|^2}{\hbar\sqrt{4\pi\lambda k_{\mathrm{B}}T}} \, \mathrm{has \ units \ of \ s^{-1}}$$

$$A \rightleftharpoons X^{\neq} \rightleftharpoons B$$

Eyring-Polanyi-Evans (1930s)

theoretical rate constant equation (from transition-state theory / activated complex theory)

$$k_f = \kappa \nu K^{\neq} = \frac{\kappa k_{\rm B} T}{h} K^{\neq}$$
 ... with transmission coefficient, κ , and vibrational frequency, ν (s⁻¹) ... and $R = N_A k_{\rm B}$... and $\frac{\kappa k_{\rm B} T}{h}$ has units of s⁻¹

$$k_f = \frac{\kappa k_{\rm B} T}{h} \exp\left(-\frac{\Delta G^{\neq}}{RT}\right) = \frac{\kappa k_{\rm B} T}{h} \exp\left(\frac{\Delta S^{\neq}}{R}\right) \exp\left(-\frac{\Delta H^{\neq}}{RT}\right) \text{ ... and so } A \text{ contains } \Delta S^{\neq}$$

... what is the largest predicted pre-exponential factor at 25 °C? (161 fs)⁻¹ = (1.61 x 10⁻¹³ s)⁻¹

proportional to

Marcus-Hush Theory

 $\frac{D + A}{A} \rightleftharpoons \frac{D^+ + A^-}{B}$

Marcus-Hush (1950s-1960s)

theoretical (semiclassical) rate constant equation

$$k_{\rm ET} = \frac{2\pi}{\hbar} |H_{\rm DA}|^2 \frac{1}{\sqrt{4\pi\lambda_{\rm AB}kT}} \exp\left(-\frac{\Delta G_{\rm AB}^{\neq}}{kT}\right)$$

quantum adiabatic electronic coupling

classical nuclear freeenergy dependence

proportional to $\exp\left(\left(\lambda_{AB} + \Delta G_{AB}^{0}\right)^{2}\right)$

... which term shall we discuss first?

... I don't know about you, but when given a choice between classical and quantum mechanical

... I go with classical... OK?

$$\Delta G_{AB}^{\neq} = \frac{\left(\lambda_{AB} + \Delta G_{AB}^{0}\right)^{2}}{4\lambda_{AB}}$$

... wait... there is a direct relation between thermodynamics (ΔG_{AB}^{o}) and kinetics (ΔG_{AB}^{\neq})?

... this makes sense since $\frac{d\overline{\mu}_A}{dz}$ drives transport, which can be written as a chemical reaction

... but that key relationship is quadratic... meaning a parabola? Huh?!?!?!

$$\frac{D + A}{A} \rightleftharpoons \frac{D^+ + A^-}{B}$$

Marcus-Hush (1950s-1960s)

theoretical (semiclassical) rate constant equation

$$k_{\rm ET} = \frac{2\pi}{\hbar} |H_{\rm DA}|^2 \frac{1}{\sqrt{4\pi\lambda_{\rm AB}kT}} \exp\left(-\frac{\Delta G_{\rm AB}^{2}}{kT}\right)$$

quantum adiabatic electronic coupling

classical nuclear freeenergy dependence

$$m{k_{ ext{ET}}} = rac{2\pi}{\hbar} |m{H_{ ext{DA}}}|^2 rac{1}{\sqrt{4\pi\lambda_{ ext{AB}}kT}} \exp\left(-rac{\left(\lambda_{ ext{AB}} + \Delta G_{ ext{AB}}^{o}
ight)^2}{4\lambda_{ ext{AB}}kT}
ight)$$
 ... what kind of function is this?
$$\Delta G_{ ext{AB}}^{\ \ \ \ \ } = rac{\left(\lambda_{ ext{AB}} + \Delta G_{ ext{AB}}^{o}
ight)^2}{4\lambda_{ ext{AB}}}$$

$$k_{\rm ET} = \frac{2\pi}{\hbar} |H_{\rm DA}|^2 \frac{1}{\sqrt{2\pi}\sqrt{2\lambda_{\rm AB}kT}} \exp\left(-\frac{1}{2} \left(\frac{\Delta G_{\rm AB}{}^{\rm o} - (-\lambda_{\rm AB})}{\sqrt{2\lambda_{\rm AB}kT}}\right)^2\right) \ ... \ \text{does this help at all?}$$

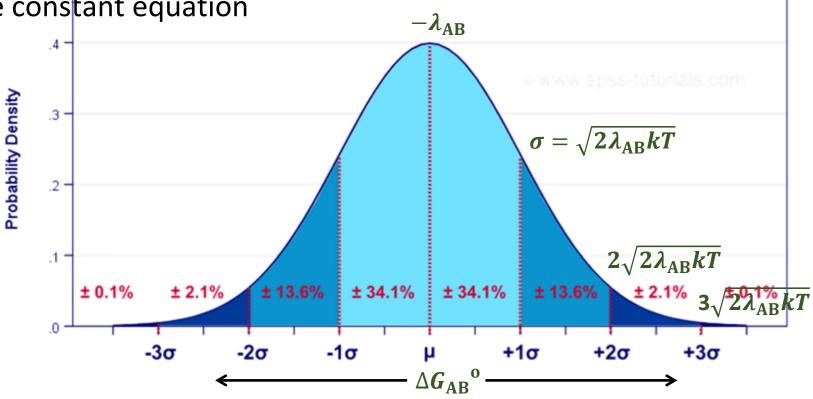
$$D + A \Longrightarrow D^+ + A^-$$

Marcus-Hush (1950s-1960s) Standard Normal Distribution A

 $\mu = 0 \mid \sigma = 1$

theoretical (semiclassical) rate constant equation

... wait... the classical component in Marcus electron-transfer theory is a **normal distribution** as a function of standard-state thermodynamic driving force (ΔG_{AB}^{0}) ?... **Yep!**

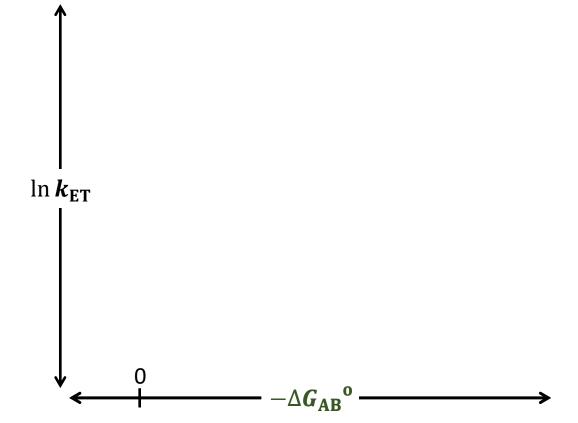


$$k_{\rm ET} = \frac{2\pi}{\hbar} |H_{\rm DA}|^2 \frac{1}{\sqrt{2\pi}\sqrt{2\lambda_{\rm AB}kT}} \exp\left(-\frac{1}{2} \left(\frac{\Delta G_{\rm AB}{}^{\rm o} - (-\lambda_{\rm AB})}{\sqrt{2\lambda_{\rm AB}kT}}\right)^2\right) \ ... \ \text{does this help at all?}$$

Marcus-Hush Theory $D + A \rightleftharpoons D^+ + A^-$

Marcus-Hush (1950s-1960s)

$$k_{\rm ET} = \frac{2\pi}{\hbar} |H_{\rm DA}|^2 \frac{1}{\sqrt{4\pi\lambda_{\rm AB}kT}} \exp\left(-\frac{\left(\lambda_{\rm AB} + \Delta G_{\rm AB}^{\rm o}\right)^2}{4\lambda_{\rm AB}kT}\right)$$



 $\frac{D + A}{A} \Longleftrightarrow \underline{D^+ + A^-}$

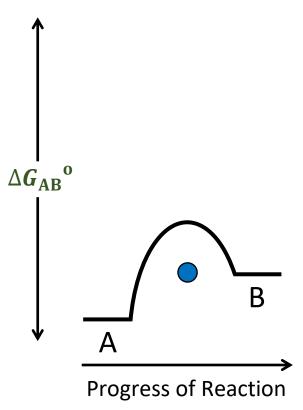
Marcus-Hush (1950s-1960s)

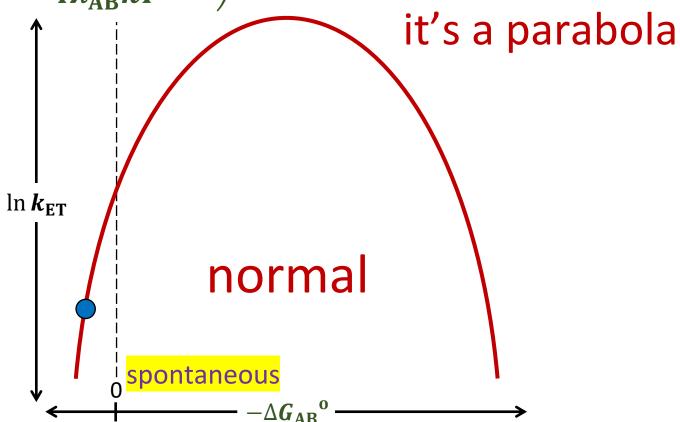
theoretical (semiclassical) rate constant equation

$$k_{\rm ET} = \frac{2\pi}{\hbar} |H_{\rm DA}|^2 \frac{1}{\sqrt{4\pi\lambda_{\rm AB}kT}} \exp\left(-\frac{\left(\lambda_{\rm AB} + \Delta G_{\rm AB}^{\ 0}\right)^2}{4\lambda_{\rm AB}kT}\right)$$

How can a thermo-dynamically unfavorable reaction proceed?

Mass Action!

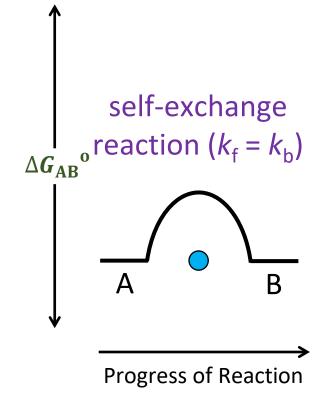


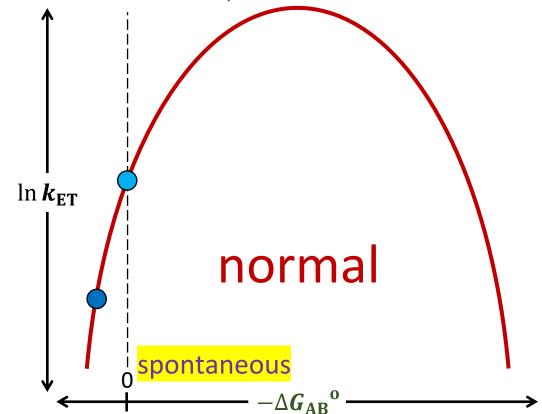


 $\frac{D + A}{A} \Longleftrightarrow \underline{D^+ + A^-}$

Marcus-Hush (1950s-1960s)

$$k_{\rm ET} = \frac{2\pi}{\hbar} |H_{\rm DA}|^2 \frac{1}{\sqrt{4\pi\lambda_{\rm AB}kT}} \exp\left(-\frac{\left(\lambda_{\rm AB} + \Delta G_{\rm AB}^{\rm o}\right)^2}{4\lambda_{\rm AB}kT}\right)$$





 $\frac{D + A}{A} \rightleftharpoons \frac{D^+ + A^-}{B}$

Marcus-Hush (1950s-1960s)

$$k_{\text{ET}} = \frac{2\pi}{\hbar} |H_{\text{DA}}|^2 \frac{1}{\sqrt{4\pi\lambda_{\text{AB}}kT}} \exp\left(-\frac{(\lambda_{\text{AB}} + \Delta G_{\text{AB}}^{\circ})^2}{4\lambda_{\text{AB}}kT}\right)$$

$$\lim_{\text{Progress of Reaction}} \ln k_{\text{ET}}$$
normal

Note that a

transition state

only lasts on the

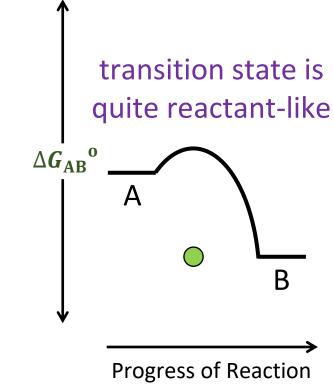
order of a bond

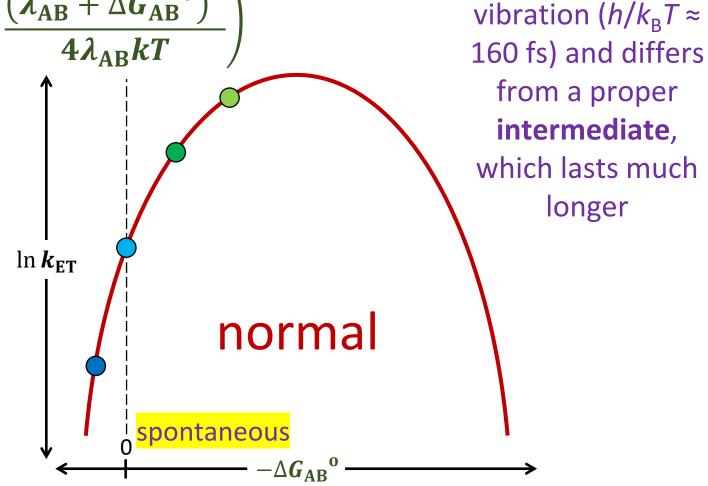
Marcus-**Hush** (1950s-1960s)

 $D + A \Longrightarrow D^+ + A^-$

theoretical (semiclassical) rate constant equation

 $k_{\rm ET} = \frac{2\pi}{\hbar} |H_{\rm DA}|^2 \frac{1}{\sqrt{4\pi\lambda_{\rm AB}kT}} \exp\left(-\frac{\left(\lambda_{\rm AB} + \Delta G_{\rm AB}^{\rm o}\right)^2}{4\lambda_{\rm AB}kT}\right)$





 $\frac{D + A}{A} \rightleftharpoons \frac{D^+ + A^-}{B}$

Marcus-Hush (1950s-1960s)

$$k_{\text{ET}} = \frac{2\pi}{\hbar} |H_{\text{DA}}|^2 \frac{1}{\sqrt{4\pi\lambda_{\text{AB}}kT}} \exp\left(-\frac{\left(\lambda_{\text{AB}} + \Delta G_{\text{AB}}^{\circ}\right)^2}{4\lambda_{\text{AB}}kT}\right)$$

$$\log_{\text{Progress of Reaction}} \log_{\text{Progress of Reacti$$

Marcus-Hush Theory $D + A \rightleftharpoons D^+ + A^-$

Marcus-Hush (1950s-1960s)

$$k_{\text{ET}} = \frac{2\pi}{\hbar} |H_{\text{DA}}|^2 \frac{1}{\sqrt{4\pi\lambda_{\text{AB}}kT}} \exp\left(-\frac{\left(\lambda_{\text{AB}} + \Delta G_{\text{AB}}^{\circ}\right)^2}{4\lambda_{\text{AB}}kT}\right)$$

$$\lim_{K \to \infty} k_{\text{ET}} = \frac{2\pi}{\hbar} |H_{\text{DA}}|^2 \frac{1}{\sqrt{4\pi\lambda_{\text{AB}}kT}} \exp\left(-\frac{\left(\lambda_{\text{AB}} + \Delta G_{\text{AB}}^{\circ}\right)^2}{4\lambda_{\text{AB}}kT}\right)$$

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 $\frac{D + A}{A} \rightleftharpoons \frac{D^+ + A^-}{B}$

Marcus-Hush (1950s-1960s)

$$k_{\rm ET} = \frac{2\pi}{\hbar} |H_{\rm DA}|^2 \frac{1}{\sqrt{4\pi\lambda_{\rm AB}kT}} \exp\left(-\frac{(\lambda_{\rm AB} + \Delta G_{\rm AB}^{\rm o})^2}{4\lambda_{\rm AB}kT}\right)$$
... this looks odd...
... but what did we expect from a $\Delta G_{\rm AB}^{\rm o}$ and $\Delta G_{\rm AB}^{\rm o}$ but what did we expect from a General Chemistry cartoon inverted inverted

Outer Reorganization Energy

Marcus-Hush (1950s-1960s)

theoretical (semiclassical) rate constant equation

$$k_{\rm ET} = \frac{2\pi}{\hbar} |H_{\rm DA}|^2 \frac{1}{\sqrt{4\pi\lambda_{\rm AB}kT}} \exp\left(-\frac{\Delta G_{\rm AB}^{\neq}}{kT}\right)$$

$$\Delta G_{AB}^{\neq} = \frac{\left(\lambda_{AB} + \Delta G_{AB}^{0}\right)^{2}}{4\lambda_{AB}}$$

... but what causes this parabolic relationship?...

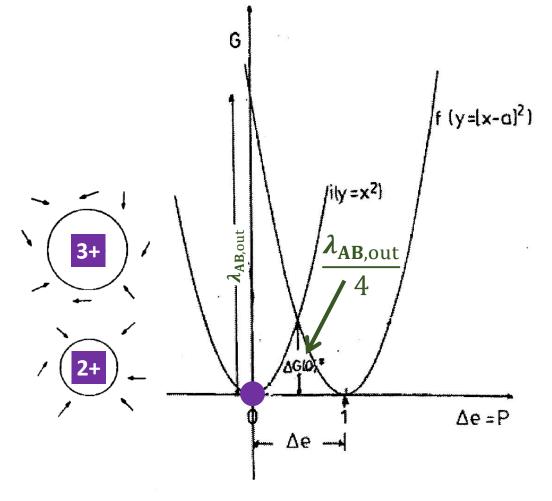
... what physical expression results in a quadratic dependence on **charge**?

... classical electrostatics... of the solvent!

... dielectric continuum model...

...
$$G_{\text{out}} \approx 1 \text{ eV}$$
 in water

... G_{out} decreases as permittivity decreases



$$\lambda_{\text{out}} = \left(\frac{1}{2r_{\text{D}}} + \frac{1}{2r_{\text{A}}} - \frac{1}{R_{\text{DA}}}\right) \left(\frac{1}{\epsilon_{\infty}} - \frac{1}{\epsilon_{\text{S}}}\right) (\Delta e)^{2}$$

Outer + Inner Reorganization Energies

Marcus-Hush (1950s-1960s)

theoretical (semiclassical) rate constant equation

This theory came about by answering the following question: For an electron-transfer event, how does one satisfy the Franck–Condon principle and the conservation of energy?

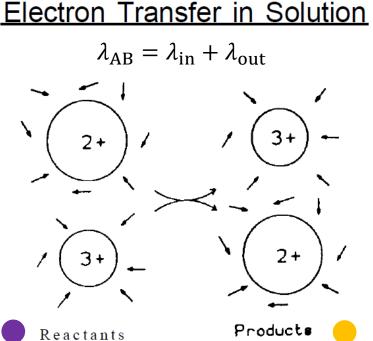
• Franck-Condon principle: Nuclei are fixed during electron-transfer between orbitals... Born-Oppenheimer

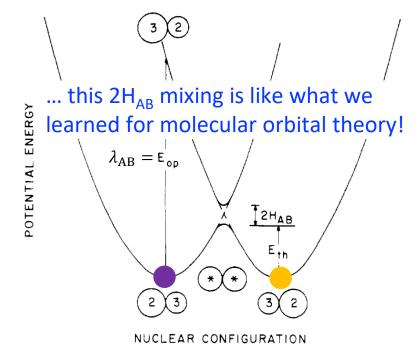
approximation is relevant

$$\lambda_i = \sum_{l} \lambda_{i,l} = \frac{1}{2} \sum_{l} f_l \left(\Delta q_{e,l} \right)^2$$

... this is for a harmonic oscillator ... with force constants, f_i ...

... and yes, it's another parabola!





Today's Critical Guiding Question

What thermodynamic term indicates whether a processes will net occur, or not, for a chemical reaction, and a single species... and uniquely for today, how does it allow us to quantitatively predict kinetic information... for real this time?

Transport as a Chemical Reaction

$$\frac{\partial c_{A,z_0}}{\partial t} = \sum_{j} R_{A,j} - \frac{\partial N_A}{\partial z}$$

... in a reaction "volume" z_0 ...

... that encompasses z_1 and z_2

$$A^-_{z_1} \rightleftharpoons A^-_{z_2}$$

$$R_{\rm A} = -k_f c_{{
m A}_{z_1}} + k_b c_{{
m A}_{z_2}} = (k_b - k_f) c_{{
m A}_{z_1}} + k_b \Delta c_{
m A}$$
 ... where $c_{{
m A}_{z_2}} = c_{{
m A}_{z_1}} + \Delta c_{
m A}$

$$\mathbf{N_{A}} = \left(-\frac{D_{A}}{RT}\frac{\partial\bar{\mu}_{A}}{\partial z}\right)\boldsymbol{c_{A}} = -\frac{D_{A}}{RT}\left(\frac{RT}{\boldsymbol{c_{A}}}\frac{d\boldsymbol{c_{A}}}{dz} + \frac{d\mu_{A}^{o}}{dz} - RT\frac{d(\ln\gamma)}{dz} - RT\frac{d(\ln\gamma)}{dz} - RT\frac{d(\ln\gamma)}{dz} - F\frac{d\varphi}{dz}\right)\boldsymbol{c_{A}}$$

$$d\boldsymbol{c_{A}} - D_{A}F\mathbf{E'} \qquad D_{A} - D_{A}F\mathbf{E'}$$

$$= -D_{A} \frac{d\mathbf{c}_{A}}{dz} + \frac{-D_{A}F\mathbf{E}'}{RT} \mathbf{c}_{A} \approx -\frac{D_{A}}{l} \Delta \mathbf{c}_{A} + \frac{-D_{A}F\mathbf{E}'}{RT} \mathbf{c}_{A}$$

... where \mathbf{E}' is an **effective force field**

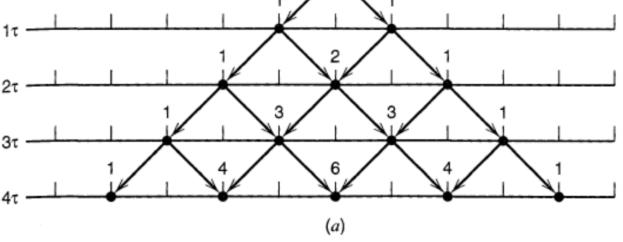
$$-\frac{\partial N_{A}}{\partial z} \approx \frac{D_{A}}{l^{2}} \Delta c_{A} + \frac{D_{A}FE'}{RTl} c_{A} = "k_{b}" \Delta c_{A} + "(k_{b} - k_{f})" c_{A}$$
(diffusion) (drift) ... what is D_{A} again?

Progress of Reaction

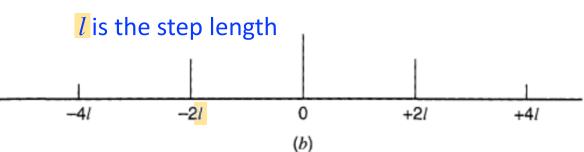
... since the reaction volume is >1 Å³ , l > 1 Å... and $D(aq) < 10^{-4}$ cm²/s... $k_{D(aq)} < 10^{12}$ s⁻¹ = (>1 ps)⁻¹

Diffusion Coefficient

m is just the number τ is the step time



root mean square (rms)
displacement
(standard deviation)



 $\overline{\Delta} = \sqrt{(2d)Dt}$, where *d* is the dimension

... and the "2" is for positive and negative directions

mean square displacement (variance) $\overline{\Delta^2} = ml^2 = \frac{t}{\tau}l^2 = 2Dt$ $\dots D = \frac{l^2}{2}$

(<u>Recall</u>: D_z is has units of cm² s⁻¹... as zz)

... and why do we care?

Figure 4.4.2 (a) Probability distribution for a one-dimensional random walk over zero to four time units. The number printed over each allowed arrival point is the number of paths to that point. (b) Bar graph showing distribution at $t = 4\tau$. At this time, probability of being at x = 0 is 6/16, at $x = \pm 2l$ is 4/16, and at $x = \pm 4l$ is 1/16.

Bard & Faulkner, Chapter 4, Figure 4.4.2, Page 147